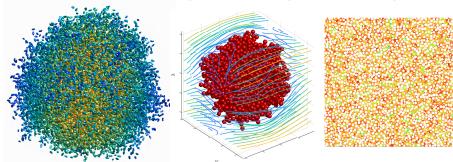
An accurate method of fundamental solutions for large and dense particle systems in Stokes flow



Anna Broms, Dept. of Mathematics, KTH, Stockholm

Joint work with

Anna-Karin Tornberg (KTH), Alex Barnett (Flatiron Institute)



Setting: tiny rigid particles in a viscous fluid

On the scale of bacteria or red blood cells (nm- μ m)

Motivated by studies of rheology, transport or diffusion processes

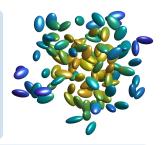
The Stokes model for the fluid:

Motivation

 A linear and viscous-dominant version of Navier-Stokes, valid in the limit of small Reynolds numbers

$$Re = UL/\nu$$

No inertia (no acceleration!)



Force balance on the particles:

 Particles interact hydrodynamically (via the fluid) and are also affected by external forces (e.g. gravity, electrostatics) Resistance Mobility Close interactions Summary 3/20

Two Stokes boundary value problems

- At a snapshot in time, the Stokes PDEs determine flow velocity u and pressure P in the fluid domain
- No-slip bc: Flow velocity at particle surface = rigid body motion of the particle
- Solution pair (u, P) relates to hydrodynamic forces and torques (= external forces and torques)

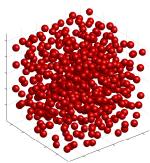
The resistance problem:

Motivation

- Given rigid body motion
- Solve for external forces and torques

The mobility problem:

- The inverse problem
- To be solved to study dynamics

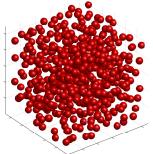


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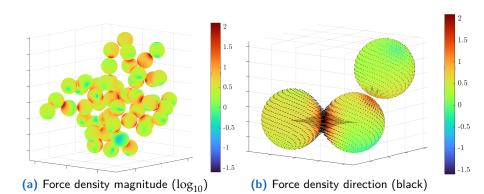
Challenges:

- Hydrodynamic interactions are long-ranged!
- Close interactions are hard to resolve!



Challenging close interactions

Relative motion results in strong *lubrication forces* for closely interacting particles and peaked force densities on their surfaces.



How?

Technique: The method of fundamental solutions (MFS)

Wish list for solver:

- Controllable accuracy, also for close interactions
- Coarse particle representations
- Scalable

Roadmap:

- Problems with weak lubrication forces: basic algorithms for
 - Resistance
 - 2 Mobility
- Problems with close interactions and strong lubrication forces:
 - 3 Image enhancement
 - 4 Preconditioning

Fundamental solutions

Motivation

The Stokes equations are given by

$$\begin{cases} \nabla P - \nabla^2 \pmb{u} = \pmb{0}, & \text{in fluid domain} \\ \nabla \cdot \pmb{u} = 0, & \text{in fluid domain} \\ \pmb{u}_{\text{bc}} = \text{rigid body motion,} & \text{on any particle surface} \\ \pmb{u} \text{ flow velocity, } P \text{ pressure.} \end{cases}$$

Stokes fundamental solutions

 The Stokeslet, S, gives velocity solution due to a point force $f = f^*\delta(x - y)$:

$$u(x) = S(x - y) f^*, \quad x \in \mathbb{R}^3 \setminus \{y\}$$

- Other fundamental solutions exist too, called the rotlet, stresslet and potential dipole
- All are singular at x = y

Basic MFS algorithm for the resistance problem

• The linear combination $u(x) = \sum_{j=1}^N \mathbb{S}(x-y_j)\lambda_j$ solves the set of PDEs, with **strategically** picked source locations y_j and **if** $\lambda_j \in \mathbb{R}^3$ can be chosen to match the boundary conditions.

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solves the set of PDEs, with strategically picked source locations y_i and if $\lambda_i \in \mathbb{R}^3$ can be chosen to match the boundary conditions.

MFS has long history:



Katsurada (1989), Götz (2003), Fairweather et al. (2005), Betcke & Barnett (2008), Alves (2009), Liu & Barnett (2016)...

Basic MFS algorithm for the resistance problem

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- Let the N source points y_j lie on a proxy-surface in the particle interior.
- Impose bc in a set of M > N target (collocation) points x_i at the particle surface.

Basic MFS algorithm for the resistance problem

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- Let the N source points y_i lie on a proxy-surface in the particle interior.
- Impose bc in a set of M > N target (collocation) points x_i at the particle surface.
- As surfaces are different, singularities in S are avoided!
- Let G be the target-from-source matrix, with $G_{ii} := \mathbb{S}(\mathbf{x}_i - \mathbf{y}_i)$, such that

$$G\lambda = \sum_{i=1}^{N} S(x - y_i) \lambda_j.$$

 The rectangular matrix G is ill-conditioned! Impose bc in the least-squares sense as

$$\min_{\lambda} \|G(x, y)\lambda - u(x)\|_2^2. \tag{1}$$

- With $G = U\Sigma V$, the pseudo-inverse is $G^{\dagger} = V\Sigma^{\dagger}U^{T}$, where singular values are truncated below a threshold to form Σ^{\dagger} .
- Solve (1), backward-stably as

$$\lambda = V \Sigma^{\dagger} \left(U^T u(x)
ight)$$
 . Order of multiplication matters!

Motivation

Stein & Barnett (2022), Lai, Kobayashi & Barnett (2015), Malhotra & Biros (2015)

• Evaluate flow in new set of points \widetilde{x} as

$$u(\widetilde{x}) = \sum_{j=1}^{N} \mathbb{S}(\widetilde{x} - y_j) \lambda_j.$$

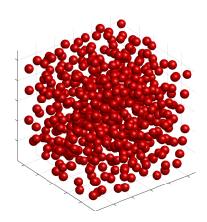
Solving for MFS coefficients (Part 2)

Assume that the system consists of many particles.

Cannot do a dense SVD globally!

Motivation

• Scales as $\mathcal{O}\left(3MP\cdot(3NP)^2\right)$ for P particles, with M>N.



Solving for MFS coefficients (Part 2)

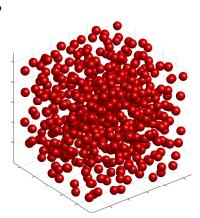
Assume that the system consists of many particles.

We construct an $\mathcal{O}(P)$ algorithm for P particles via

GMRES

Motivation

- A fast summation technique (FMM)
- Use of the block structure of the target-from-source matrix G
- One-body preconditioning:
 Only a small SVD needed to create well-conditioned system



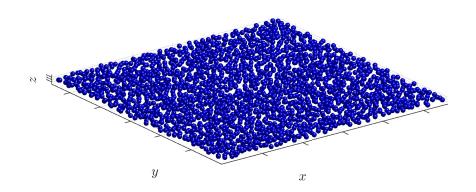
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 Motivation
 Resistance
 Mobility
 Close interactions
 Summary

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Large scale example

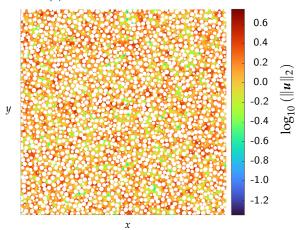
- Example: Dilute suspension of 2000 spheres in a layer.
- No-slip bc with randomly sampled angular and translational velocity on every sphere.



Large scale example

Motivation

• Max relative residual 10^{-3} , with target tolerance 10^{-3} . Residual = difference from boundary condition in new set of points on the surface of every particle.



The resistance problem (1 particle):

Solve

Motivation

 $G\lambda \approx K_M U$

for λ , with $K_M U$ the **known** rigid body motion at the surface.

• Compute forces and torques in post-processing step:

$$f = \sum_{j=1}^{N} \lambda_j, t = \sum_{j=1}^{N} (y_j - c) \times \lambda_j$$
 (*)

The mobility problem (1 particle):

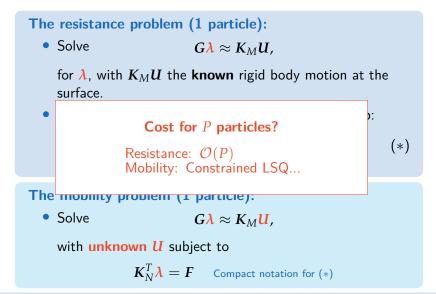
Solve

$$G\lambda \approx K_M U$$

with unknown U subject to

$$K_N^T \lambda = F$$
 Compact notation for $(*)$

Mobility formulation



Ideas:

Motivation

- **1** Create source vector λ that automatically satisfies constraints:
- 2 Split $\lambda = (I L)\hat{\lambda} + \lambda_0$, so that $K_N^T \lambda_0 = F(\lambda_0 \text{ satisfies the constraints})$ and $(I L)\hat{\lambda} \in \mathcal{N}(K_N^T)$ via projection matrix L

Proposed representation:

$$u(x) = \sum_{j=1}^{N} S(x - y_j)(I - L) \hat{\lambda}_j + \sum_{j=1}^{N} S(x - y_j) \lambda_{0,j}$$

3 Impose relation between source vector and unknown velocities: $\mathbf{U} = -K_N^T \hat{\lambda}$ (utilize unused subspace)

System to solve without explicit constraints (1 particle):

$$\left(G(I-L)+K_{M}K_{N}^{T}\right)\hat{\boldsymbol{\lambda}}=-G\boldsymbol{\lambda}_{0}$$

4 Apply one-body preconditioning

Example: 10000 ellipsoids

Motivation

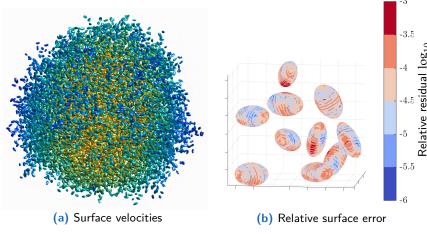


Figure: Minimum separation $\delta_{\min} = 0.2$, semi-axes: 0.4,0.6,1

Example: 10000 ellipsoids

Motivation

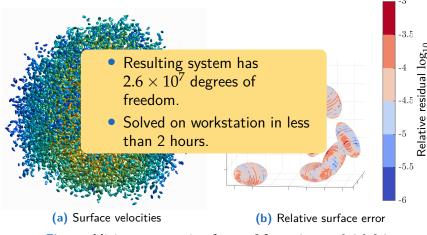


Figure: Minimum separation $\delta_{\min} = 0.2$, semi-axes: 0.4,0.6,1

Accuracy in MFS with coarse grids: Image points

Idea: Introduce extra sources at image points, with a combination of different types of fundamental solutions at each point.

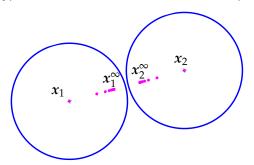




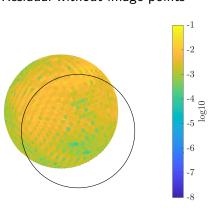
Image systems: Lord Kelvin (1853), Cheng & Greengard (1998), Cheng (2000).



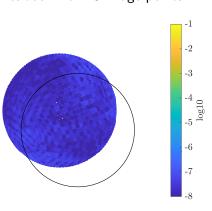
Fundamental solutions for reflection in a plane: Blake (1971), Greengard, Gimbutas & Veerapaneni (2015), Yan & Shelley (2018), Fundamental solutions for reflection of point in a sphere: Wróbel, Cortez, Varela & Fauci (2016), Maul & Kim (1996)

Example: two spheres with images, separation 10^{-2}

Residual without image points



Residual with 10 image points



Example: adaptive use of images

- Adaptive number of image points per close particle pair
- Each sphere has random translational and angular velocity
- Target accuracy of 10^{-3} reached everywhere

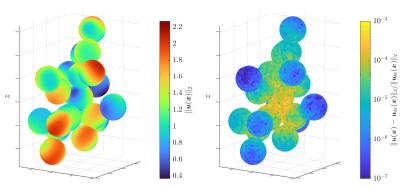


Figure: Surface velocities.

Figure: Relative residual at particle surfaces.

Preconditioning for close contacts

One-body preconditioning

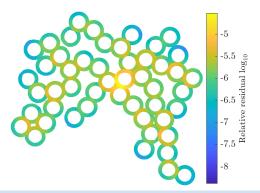
- Works well if lubrication forces are weak!
- Slow GMRES convergence if lubrication forces are strong...
- True in both 2D and 3D

New idea:

- Use a pair corrected basis inspired by
 - Cheng & Greengard (1998)
- Flavour of a fast direct solver: solve a BVP with a fine grid only locally for each close pair of particles
- Global solve only needs coarse grid

Preconditioning mobility example in 2D

- Each unit circle $\delta = 10^{-3}$ from at least one neighbour, random net forces and torques (very challenging)
- 297 iterations with one-body preconditioning, 57 with the new idea to reach GMRES tol 10^{-6}
- 60 points per particle in coarse grid



Summary

Motivation

The method of fundamental solutions

- Singularity free technique
- Additional sources at approximate image locations improve accuracy for closely interacting particles
- Controllable accuracy obtained for all separations $> 10^{-3}$
- Well-conditioned solvers for both resistance and mobility (work in progress for close contacts – talk to me!)
- Coarse representations and scalable to many particles



Accurate Stokes flow for closely interacting spheres using lubrication-adapted image systems, A. Broms, A. H. Barnett and A.-K. Tornberg, JCP, 113636 (2025)



A method of fundamental solutions for large-scale 3D elastance and mobility problems, A. Broms, A. H. Barnett and A.-K. Tornberg, Preprint (2024) – on arXiv, submitted to ACOM, under review

- Thank you! - annabrom@kth.se