

Spectrally accurate direct solution of variable-media scattering problems via impedance maps

Banff 12/10/13

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Collaborators: Adrianna Gillman (Dartmouth), Gunnar Martinsson (Boulder)

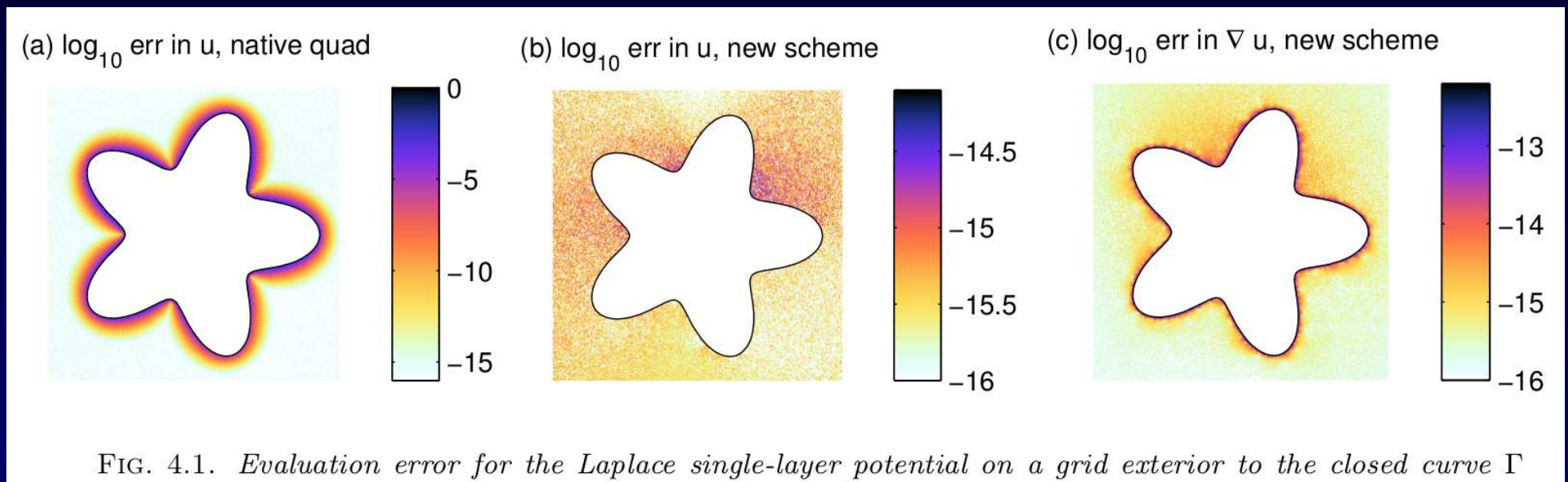


First a teaser of projects that are *not* the topic of the talk...

I. Stokes 2D close-evaluation quadratures

(w/ S. Veerapaneni, B. Wu)

- global & panel-based schemes, built on Laplace potentials & gradients
- extend Laplace global (Helsing–Ojala '08) to single-layer case...



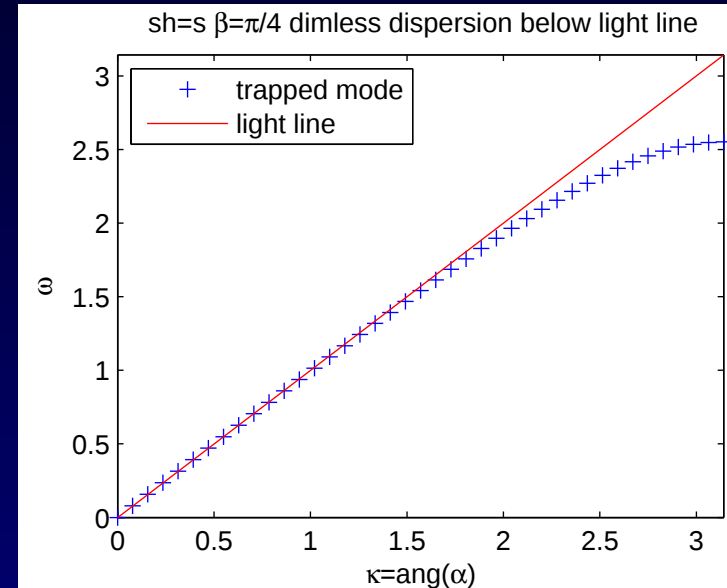
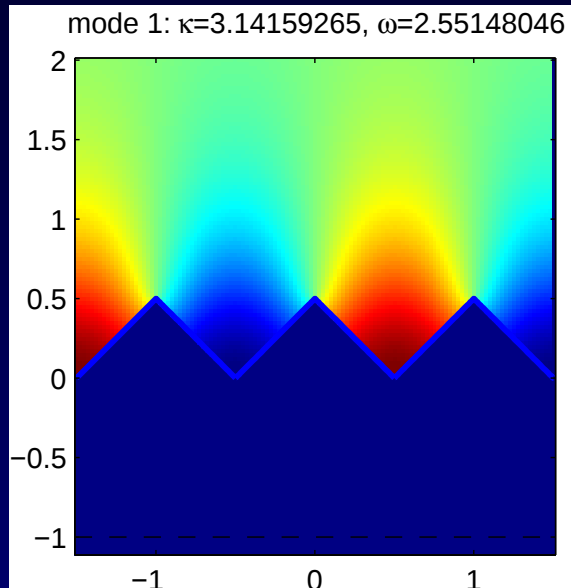
- Fix round-off instability in usual 1st derivative of barycentric form

II. Acoustic surface modes of periodic structures

(w/ E. Heller)

Footsteps on staircases at Mayan pyramid sound like *raindrops*! Why?

- each freq. travels at own group velocity $\rightarrow$ impulse becomes *chirp*

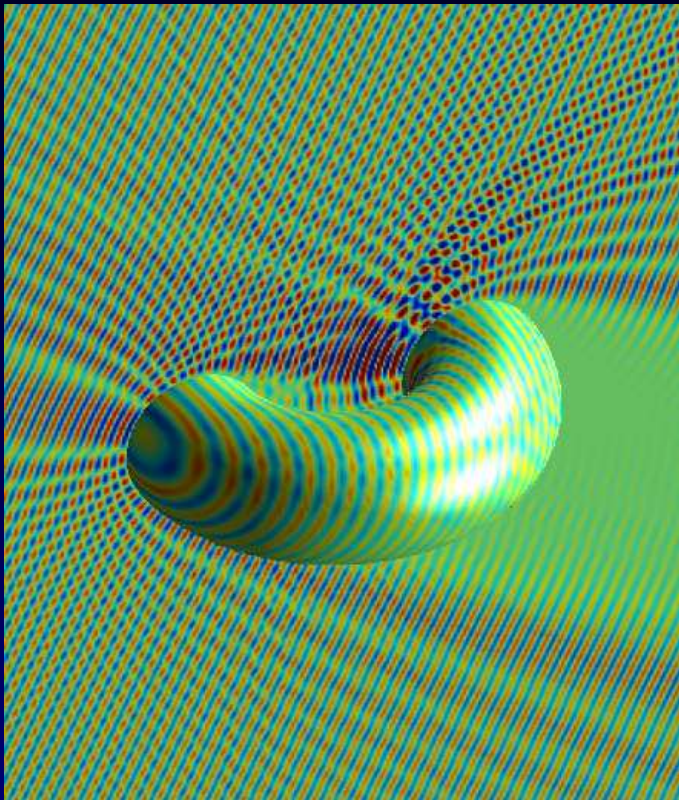


- integral formulation, periodizing via least-squares collocation on a box
- dispersion curves to spectral accuracy, via Boyd's analytic rootfinding

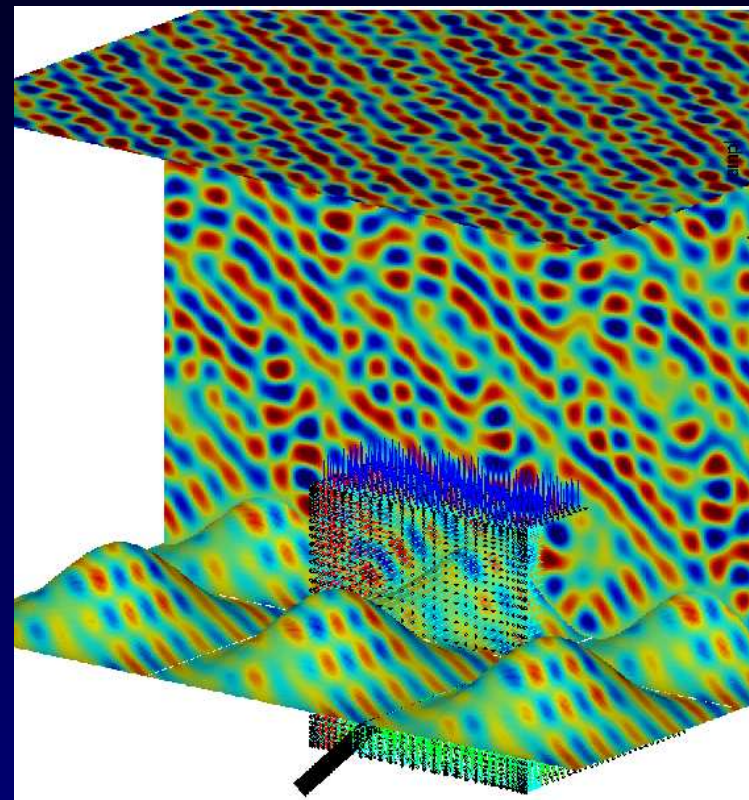
III. 3D Helmholtz QBX quadrature & periodizing

(w/ Z. Gimbutas and L. Greengard)

Sound-soft acoustic scatt: QBX setup $\approx$ 50 bare FMM GMRES iters.



30λ across, 5 digits, 2 hrs

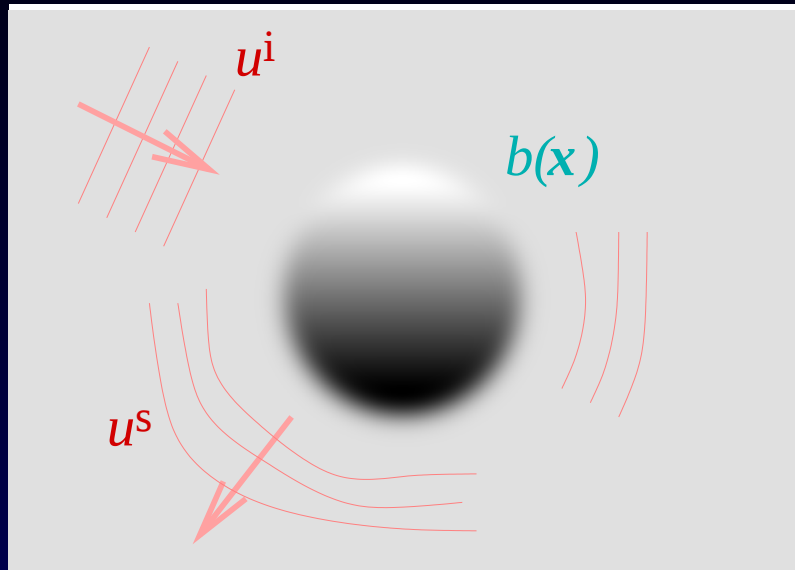


$8\lambda \times 8\lambda$ periodicity, 6 digits, 7 min

- periodizing via least-squares collocation, robust at Wood's anomalies

Please let's chat about them! Okay, now to the main talk...

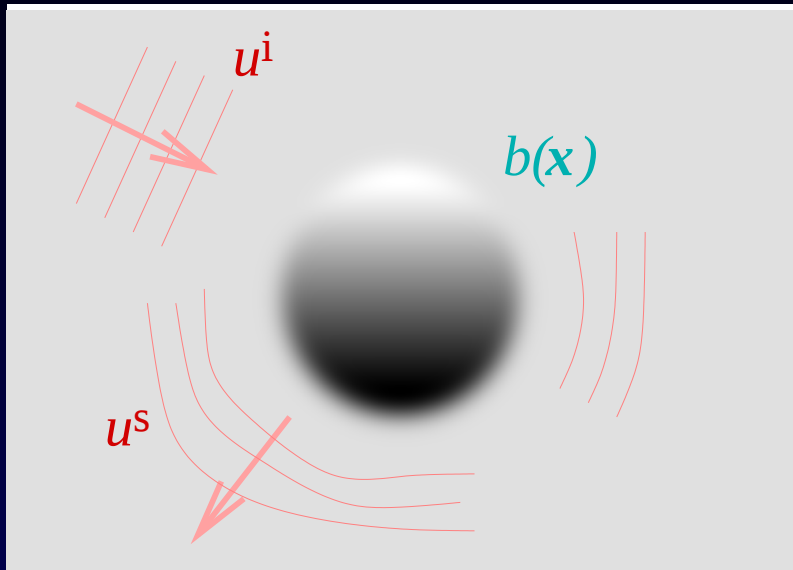
Variable-medium freq-domain scattering problem



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$$(\Delta + \kappa^2)u^i = 0 \quad \text{in } \mathbb{R}^2$$

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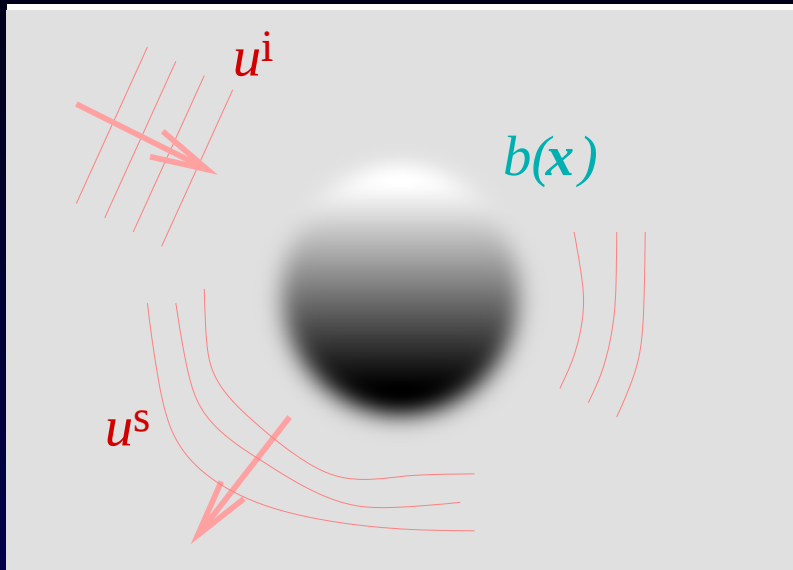
↖ scattering potential b

$\sqrt{1 - b}$ 'slowness' (refractive index)

Scattered wave u^s radiative: $\frac{\partial u^s}{\partial r} - i\kappa u^s = o(r^{-1/2})$, $r := |\mathbf{x}| \rightarrow \infty$

If $1 - b > 0$ everywhere, $\exists$ unique solution $\forall \kappa > 0$

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- Applications: underwater acoustics, seismic imaging, ultrasound, quantum chemistry & physics, optics in graded index, metamaterials

Want: smooth b , high κ , multiple incident angles

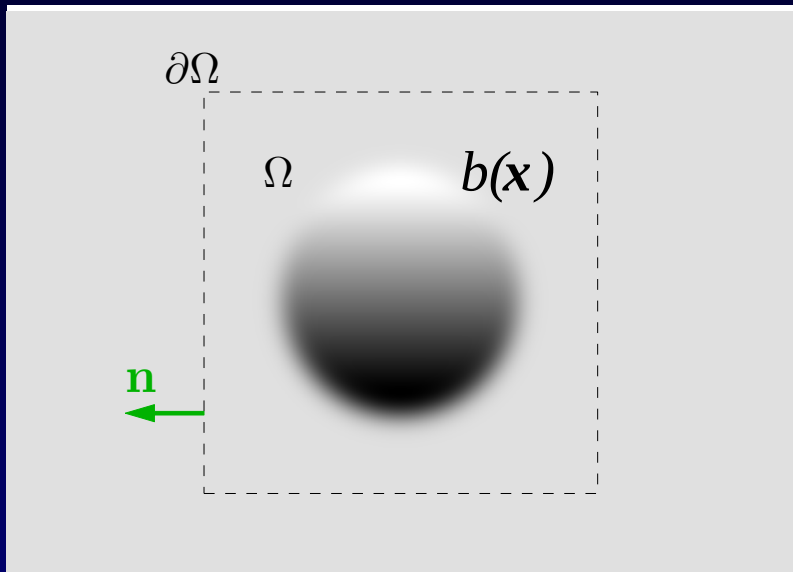
Coupling on enclosing box

Std: PDE for scatt. wave: $(\Delta + \kappa^2)u^s - \kappa^2 b u^s = \kappa^2 b u^i \quad \leftarrow \text{inhomog.}$
gives Lippmann-Schwinger volume IE...

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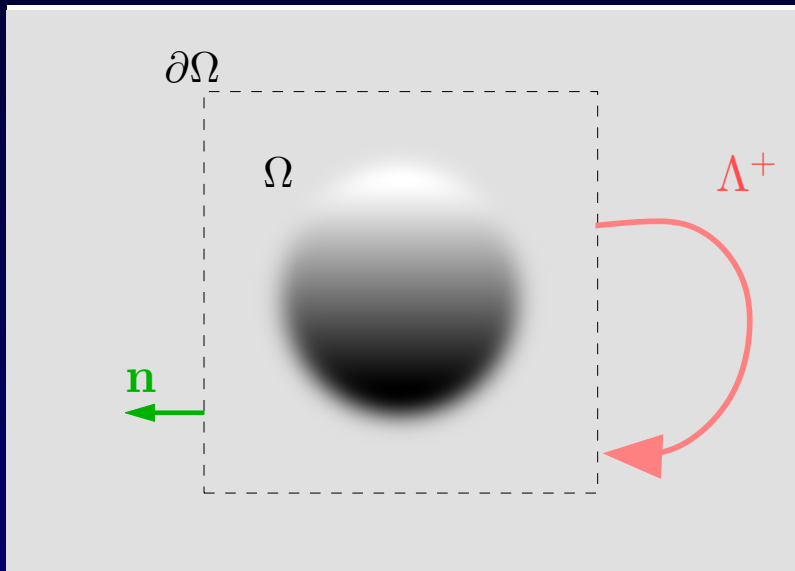
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Λ^+ : exterior Dirichlet-to-Neumann map

$$\Lambda^+ u^s|_{\partial\Omega} = \frac{\partial u^s}{\partial n} =: u_n^s$$

u^s radiative, $(\Delta + \kappa^2)u^s = 0$ in $\mathbb{R}^2 \setminus \overline{\Omega}$

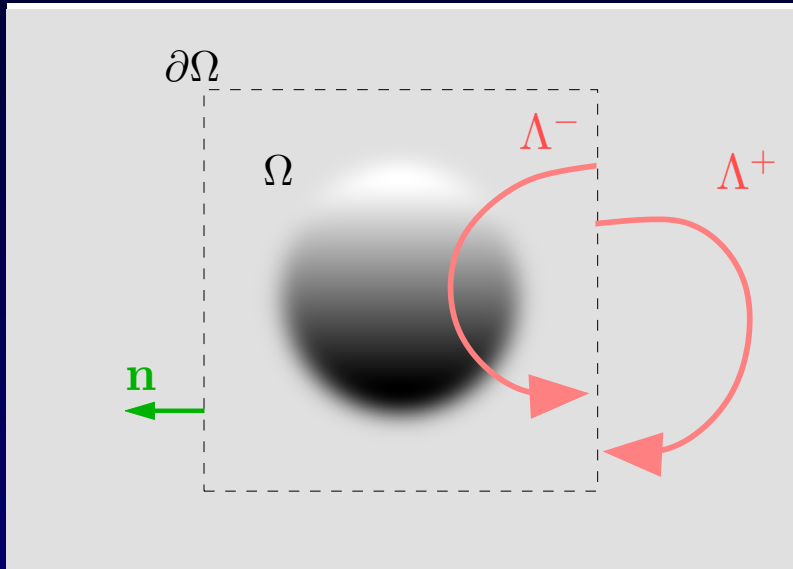
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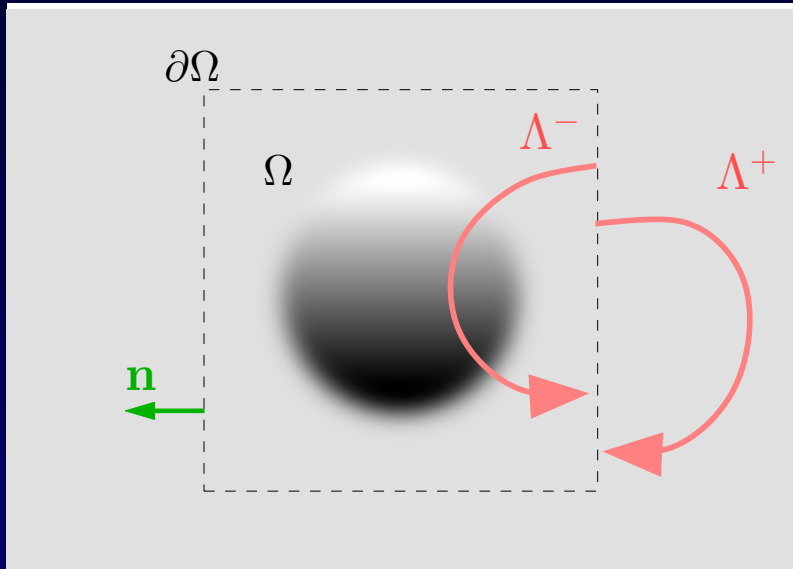
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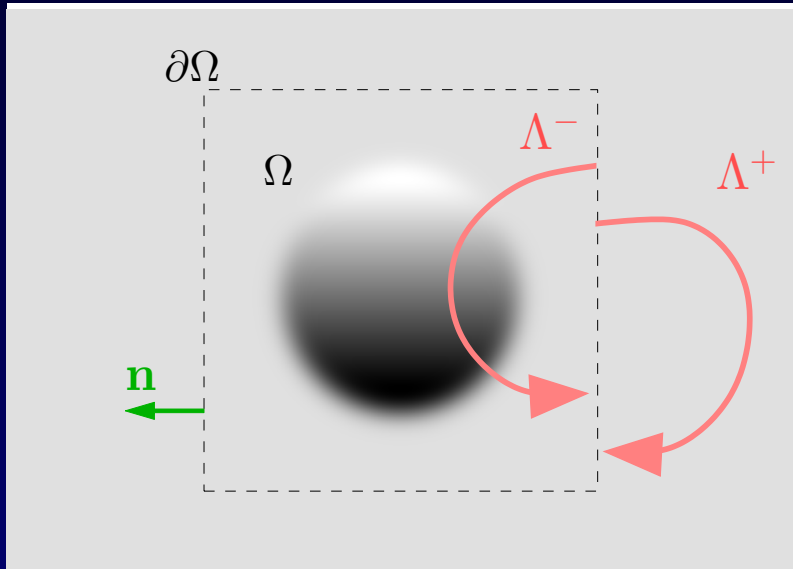
$$\Lambda^-(u^i + u^s)|_{\partial\Omega} = u_n^i + u_n^s = u_n^i + \Lambda^+ u^s|_{\partial\Omega}$$

$$\Rightarrow (\Lambda^- - \Lambda^+)u^s|_{\partial\Omega} = u_n^i - \Lambda^+ u^i|_{\partial\Omega} \quad (\text{Kirsch-Monk '94, FEM coupling})$$

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$\nwarrow$ both order +1, signs *add* $\rightarrow$ ill-cond. $\rightarrow$ inversion would lose digits

2nd kind scattering formulation

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Recall Helmholtz boundary integral operators on $\partial\Omega$:

single-layer $(S\phi)(\mathbf{x}) := \int_{\partial\Omega} \frac{i}{4} H_0^{(1)}(\kappa|\mathbf{x} - \mathbf{y}|) \phi(\mathbf{y}) ds_{\mathbf{y}}$

double-layer $(D\phi)(\mathbf{x}) := \int_{\partial\Omega} \frac{i}{4} \frac{\partial}{\partial n_{\mathbf{y}}} H_0^{(1)}(\kappa|\mathbf{x} - \mathbf{y}|) \phi(\mathbf{y}) ds_{\mathbf{y}}$

ext. Green's rep. $u|_{\partial\Omega} = (D + \frac{1}{2})u|_{\partial\Omega} - Su_n, \quad \text{so} \quad \Lambda^+ = S^{-1}(D - \frac{1}{2})$

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$$\text{Left-regularize } (*) \text{ by } S: \quad (\tfrac{1}{2} - D + S\Lambda^-)u^s|_{\partial\Omega} = S(u_n^i - \Lambda^-u^i|_{\partial\Omega})$$

Thm. Let $\partial\Omega$ be Lipschitz, $b(\mathbf{x})$ bnded, $0 < \kappa \neq \kappa_j$.

Then $A := \tfrac{1}{2} - D + S\Lambda^- = I + \text{compact}$ above BIE is 2nd kind

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Pf. let $\mathcal{P} : L^2(\partial\Omega) \rightarrow L^2(\Omega)$ be interior solution operator: $u|_{\Omega} = \mathcal{P}u|_{\partial\Omega}$

$\mathcal{V} : L^2(\Omega) \rightarrow H^{1/2}(\partial\Omega)$ be volume potential $(\mathcal{V}\phi)(\mathbf{x}) = \int_{\Omega} \frac{i}{4} H_0^{(1)}(\kappa|\mathbf{x} - \mathbf{y}|) \phi(\mathbf{y}) d\mathbf{y}$

$$S\Lambda^-u|_{\partial\Omega} = Su_n = (\tfrac{1}{2} + D)u|_{\partial\Omega} + \kappa^2 \mathcal{V}(b \cdot \mathcal{P}u|_{\partial\Omega}) \quad \text{G's 3rd id, } \forall u \in L^2(\partial\Omega)$$

$\mathcal{V}b\mathcal{P}$ cpt by Sobolev imbedding $\nearrow$ (Lipschitz bdry: McLean '00)

□

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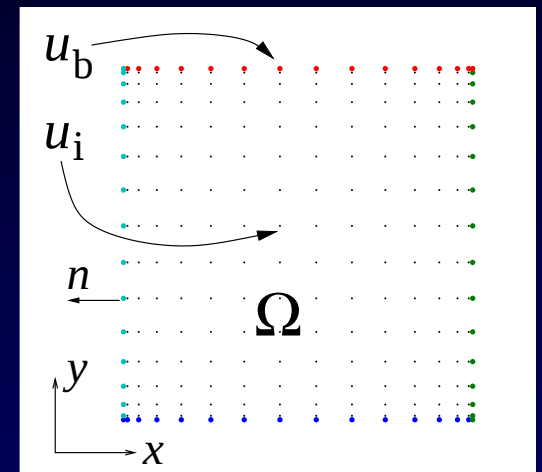
p^2 unknowns, vector $\mathbf{u} = \begin{bmatrix} u_b \\ u_i \end{bmatrix}$ $\leftarrow$ boundary $\leftarrow$ interior

$p^2 \times p^2$ spectral diff. matrices D_x, D_y act on $\mathbf{u}$

matrix $H = D_x^2 + D_y^2 + \kappa^2 \text{diag}\{1 - b(\mathbf{x}_j)\}$

$H_{i,:}$ = rows of H imposing PDE at interior nodes

system $B\mathbf{u} := \begin{bmatrix} I & 0 \\ -H_{i,:} \end{bmatrix} \begin{bmatrix} u_b \\ u_i \end{bmatrix} = \begin{bmatrix} \mathbf{f} \\ 0 \end{bmatrix}$ $\leftarrow$ Dirichlet BC (trivial) $\leftarrow$ PDE at interior nodes



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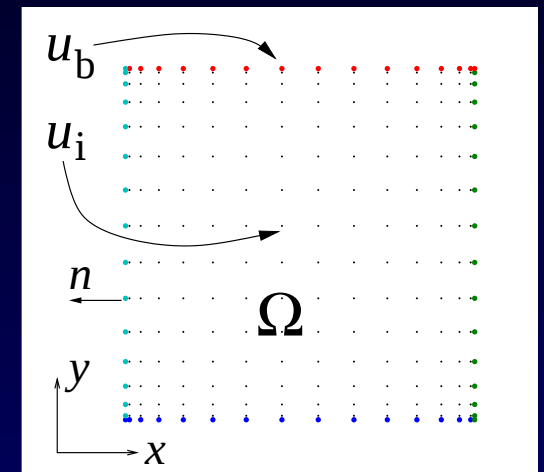
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solution matrix X , maps $\mathbf{f}$ to $\mathbf{u}$: solve $BX = \begin{bmatrix} I \\ 0 \end{bmatrix}$ dense, $O(p^6)$

Get approx. Λ^- : X left-mult. by rows of D_x, D_y to take normal deriv.

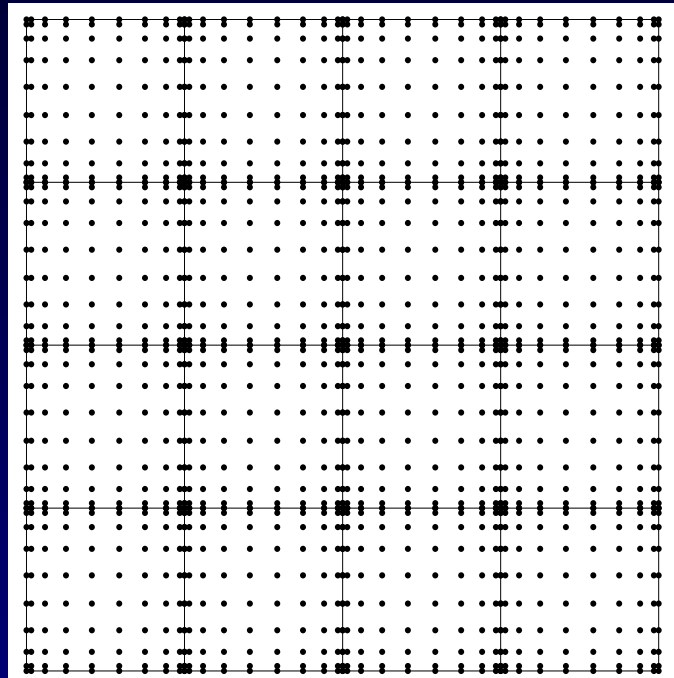
How build Λ^- for “large” scatterer?

One-box dense scheme was $O(p^6)$ terrible if need more resolution

Partition top-level box Ω into 2^L -by- 2^L small square “leaf” boxes

Give each box p^2 spectral product nodes as before total unknowns = N

Here's $L = 2$:



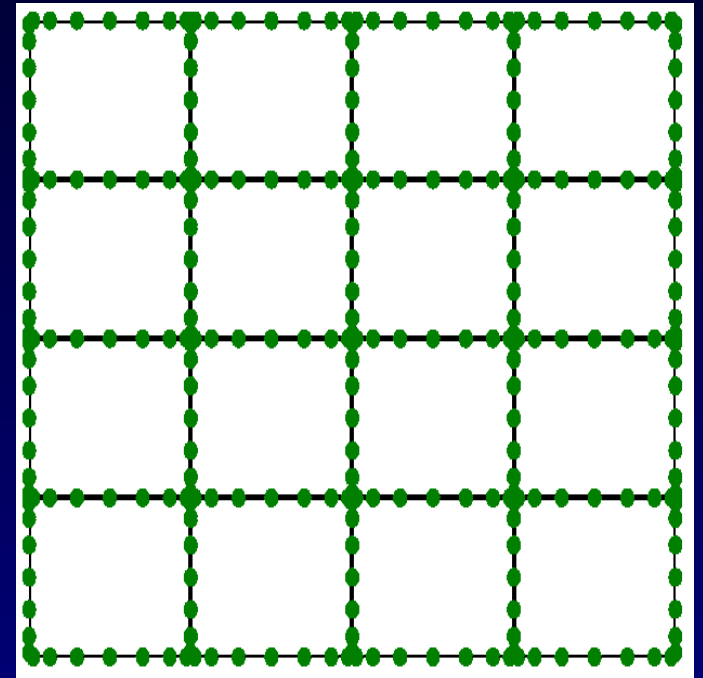
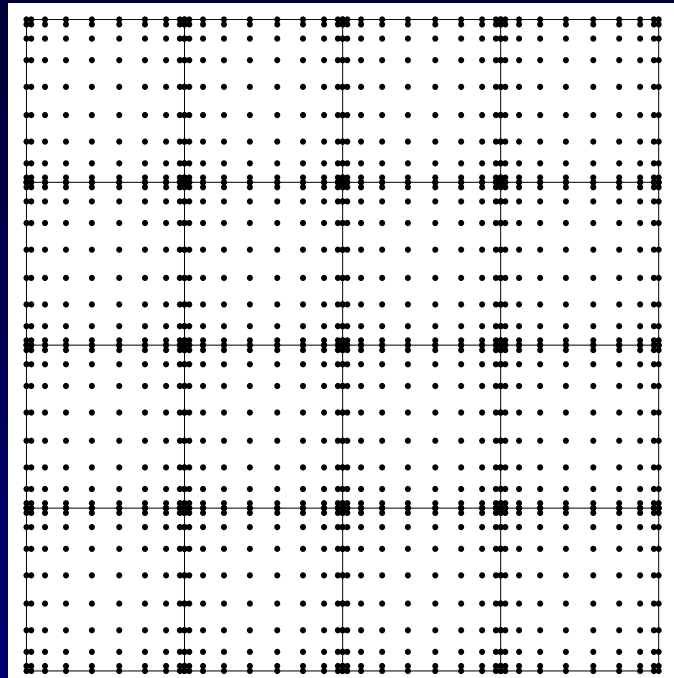
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Step 1: build Λ^- for each box separately

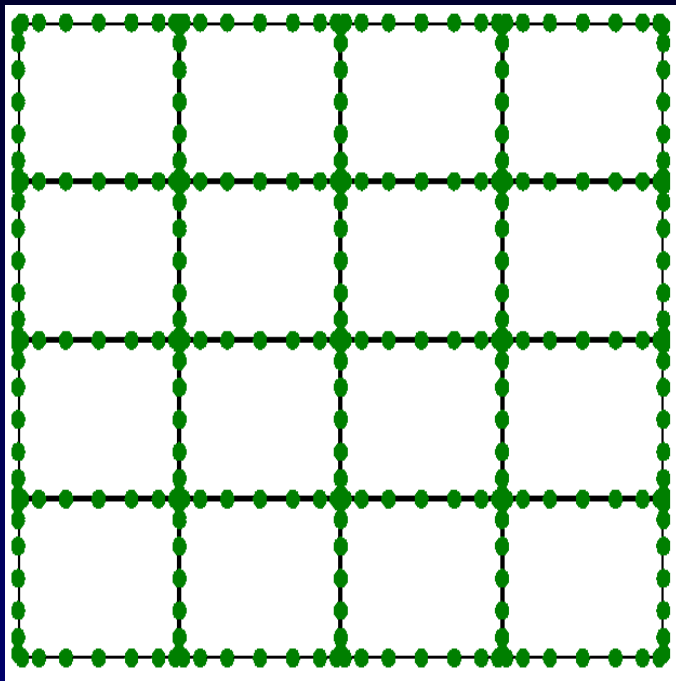
fix $p \approx 16$ effort $O(N)$

Merging leaf DtN operators

Can eliminate unknowns on common edge of two neighboring boxes

Creates Λ^- matrix for union of the boxes

dense inverse, $O(p^3)$

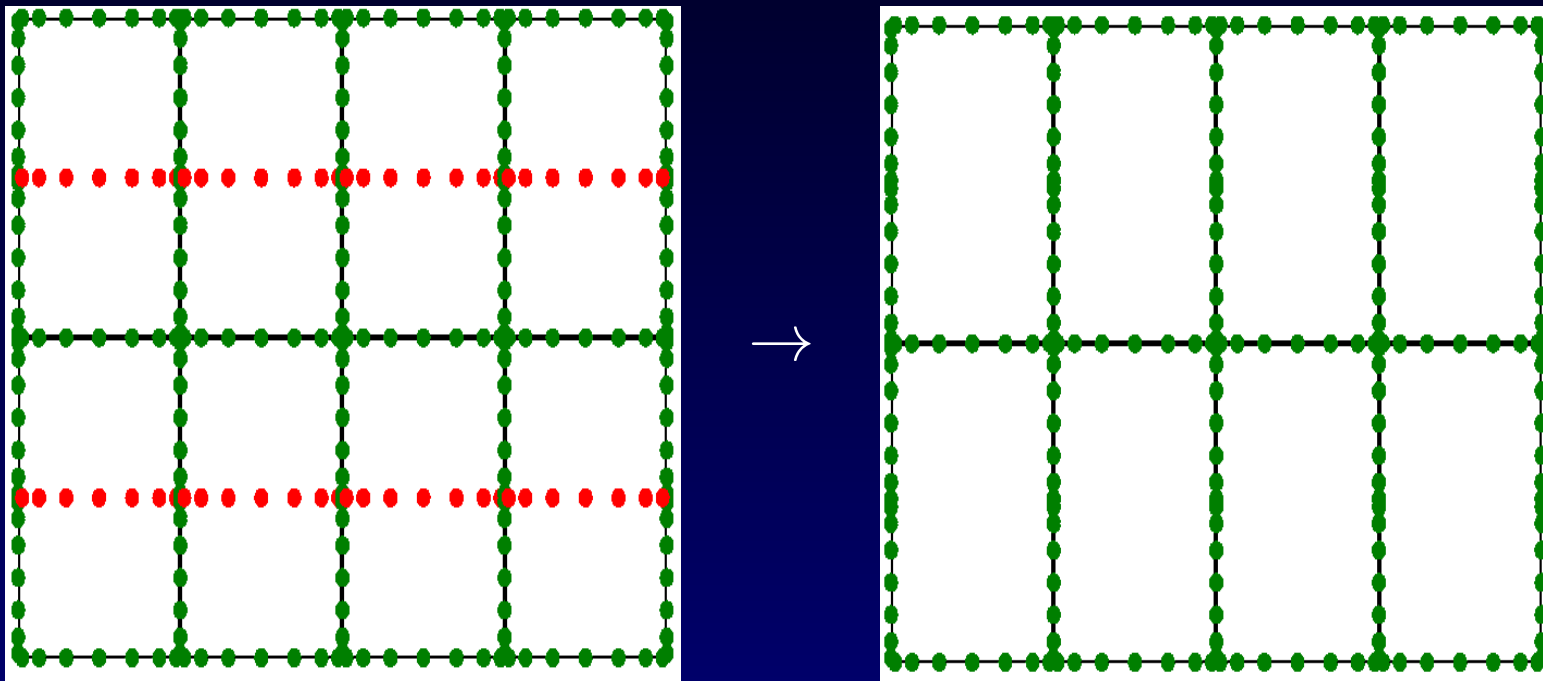


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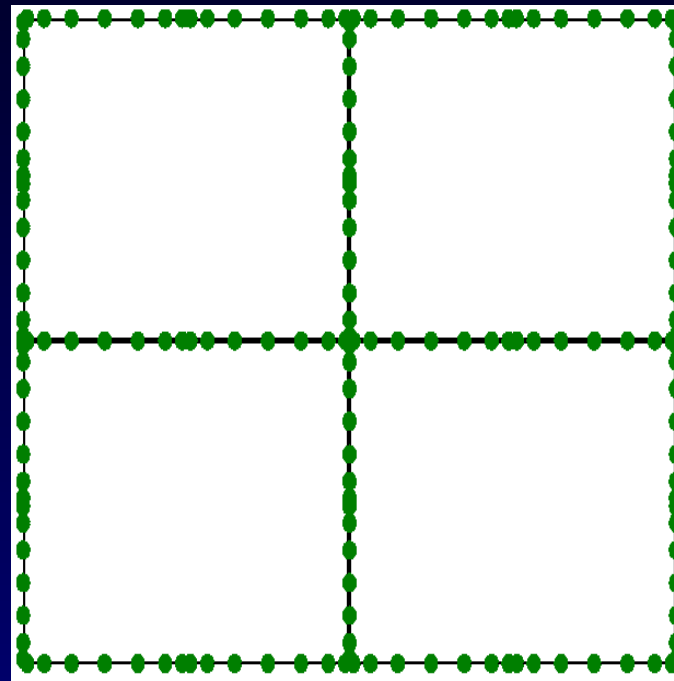
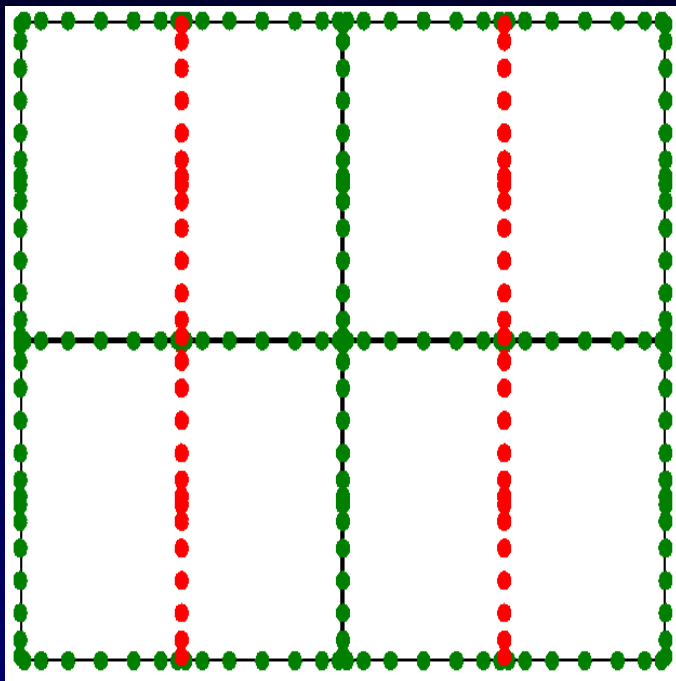
Now repeat: this is *upwards sweep* of the quad-tree...

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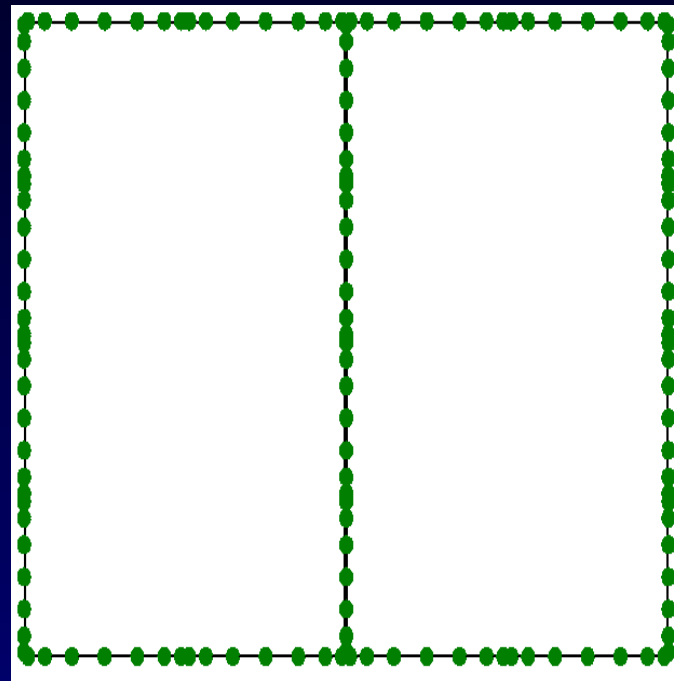
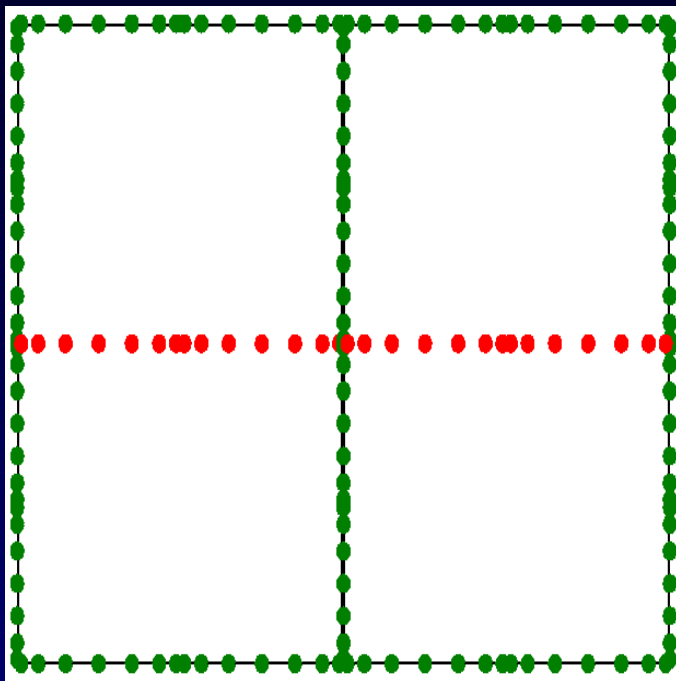


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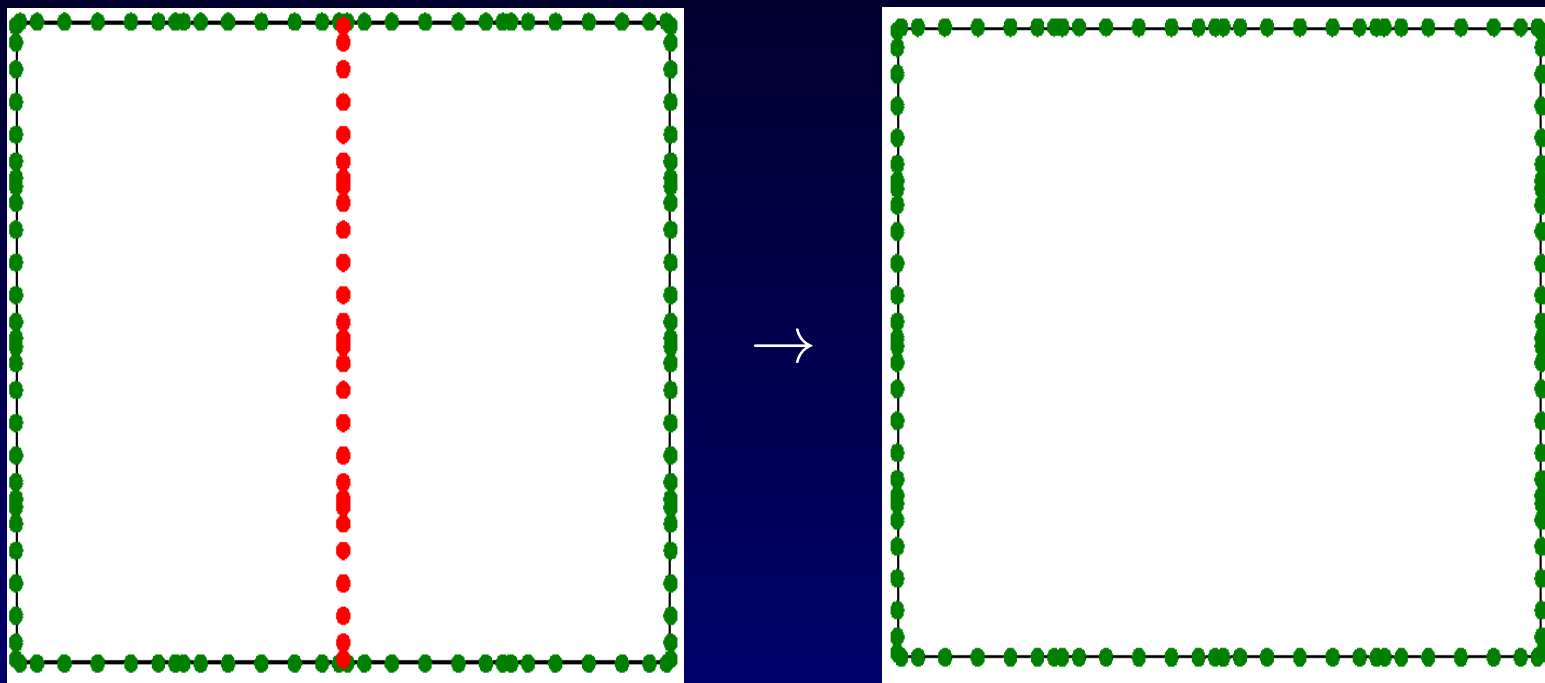


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Top-level merge is $\frac{1}{2}$ total effort: $O(N^{1/2})$ dense inverse effort $O(N^{3/2})$

We now have matrix approx. to Λ^- on edge nodes of entire box

Details: actually we use Gauss–Legendre nodes; no corner refinement (soln. u smooth)

Problem: interior DtN resonances

Recall: DtN does not exist at discrete set of Dirichlet eigenfreqs. $\kappa = \kappa_j$

$\Rightarrow$ numerically **not robust**: DtN norm blows up near κ_j , lose digits

Every leaf *and merged box* has own (unpredictable, b -dependent) eigenfreqs!

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Remedy, use $R : L^2(\partial\Omega) \rightarrow L^2(\partial\Omega)$ impedance-to-impedance map

fix η real, let $\begin{cases} f := u_n + i\eta u|_{\partial\Omega} & \text{"incoming"} \\ g := u_n - i\eta u|_{\partial\Omega} & \text{"outgoing"} \end{cases}$ then, $Rf = g$

classical facts: R exists for all κ . R is *unitary*.

Instead build ItI for all leaves upper block of B now: D_x, D_y rows + $[i\eta I \quad 0]$

Hierarchical ItI merge operations: max cond. # now 20 not 2×10^5

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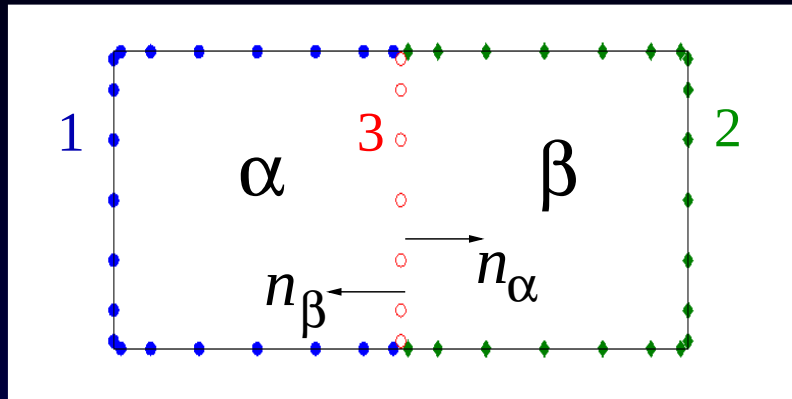
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Then convert top-level ItI to DtN:

Cayley transform $\Lambda^- = -i\eta(R - I)^{-1}(R + I)$ $O(N^{3/2})$

ItI merge step for two boxes



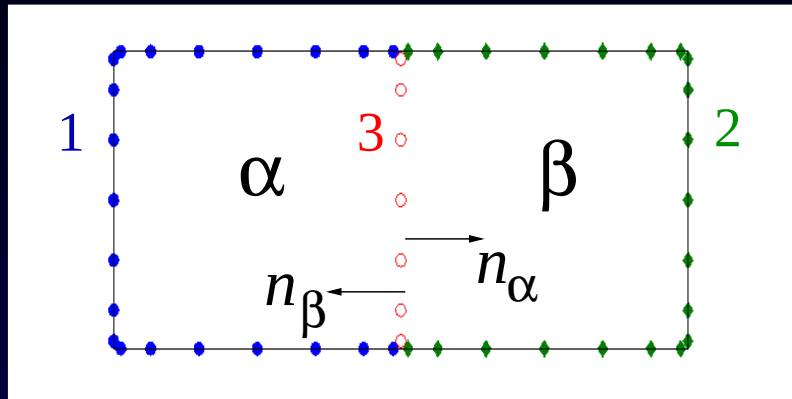
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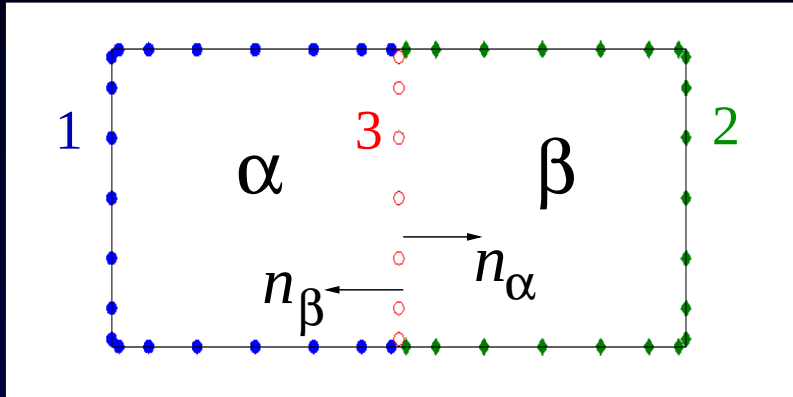
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Write edge 3 part of ItI relation $Rf = g$ in each box:

$$\begin{aligned}
 R_{31}^{\alpha} f_1 + R_{33}^{\alpha} f_3^{\alpha} &= g_3^{\alpha} &= -f_3^{\beta} & \text{by matching, normals opposed (2)} \\
 R_{32}^{\beta} f_2 + R_{33}^{\beta} f_3^{\beta} &= g_3^{\beta} &= -f_3^{\alpha}
 \end{aligned}$$

ItI merge step for two boxes



1 = nodes only on α

2 = nodes only on β

3 = shared nodes

u, u_n match on edge 3

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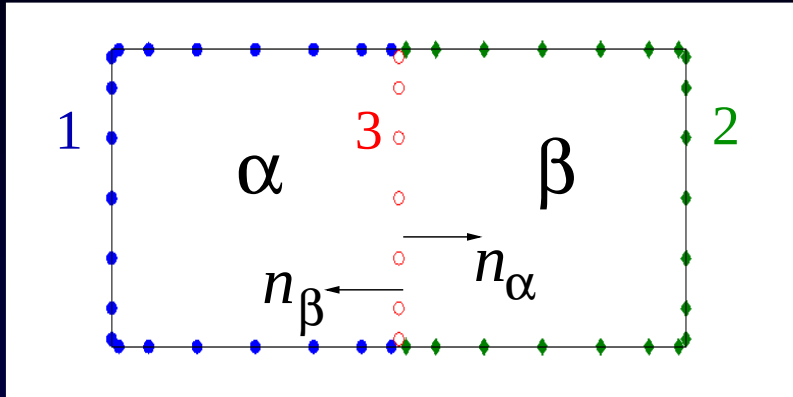
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$$\text{Elim. } f_3^{\alpha}: \quad f_3^{\beta} = F^{\beta} \begin{bmatrix} f_1 \\ f_2 \end{bmatrix}, \quad F^{\beta} := (1 - R_{33}^{\alpha} R_{33}^{\beta})^{-1} [-R_{31}^{\alpha} \quad R_{33}^{\alpha} R_{32}^{\beta}]$$

$$\text{also } F^{\alpha} := [0 \quad -R_{32}^{\beta}] - R_{33}^{\beta} F^{\beta},$$

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$$\text{also } F^{\alpha} := [0 \quad -R_{32}^{\beta}] - R_{33}^{\beta} F^{\beta}, \quad \text{then } R = \begin{bmatrix} R_{11}^{\alpha} & 0 \\ 0 & R_{22}^{\beta} \end{bmatrix} + \begin{bmatrix} R_{13}^{\alpha} F^{\alpha} \\ R_{23}^{\beta} F^{\beta} \end{bmatrix}$$

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Pre-computation:

- 1) Nyström discretization of S , D at $n = O(N^{1/2})$ nodes on $\partial\Omega$
e.g. generalized Gaussian quadrature, 6 levels dyadic corner-refinement

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- 6) eval. u inside Ω via reversing all merges, traverse down the tree
matvecs with soln. matrices F^α , etc, created during merges $O(N)$

Relation to prior work

Existing methods for variable-medium scattering:

- time-domain (FDTD): low-order, dispersion errors, long settling time
- freq-domain: low κ or low accuracy, e.g. . .
 - FEM-BIE coupling, low κ , preconditioned BIE, 3 digits (Kirsch–Monk '94)
 - 4th-order FD + PML, UMFPACK, low κ (Britt–Tsynkov–Turkel '11)
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Our influences:

- multi-domain spectral collocation (Orszag, Patera, '80s; Hesthaven, Pfeiffer)
- direct solvers for block-sparse: nested dissection, multifrontal (George, Duff, '70s)
- $O(N)$ direct solvers (Martinsson–Rokhlin '05, Xia–Chandrasekaran–Gu–Li '09)
- $O(N^{3/2})$ direct solver using both u and u_n data (Chen '02, '13)
- DtN version (Martinsson '12); ItI makes robust for oscillatory PDE

Results: time and memory

8-core Xeon, MATLAB, not optim!

t_{pre} precomputation: build DtN, build A , dense A^{-1}
 $t_{\text{new ext}}$ compute $u^s|_{\partial\Omega}$, u_n^s for new incident wave (ext. / far field)
 $t_{\text{new int}}$ propagate u from bdry down to all leaf interior nodes

$L = \# \text{ tree levels}$ $N = \# \text{ unknowns}$ $n = \# \text{ BIE unknowns}$

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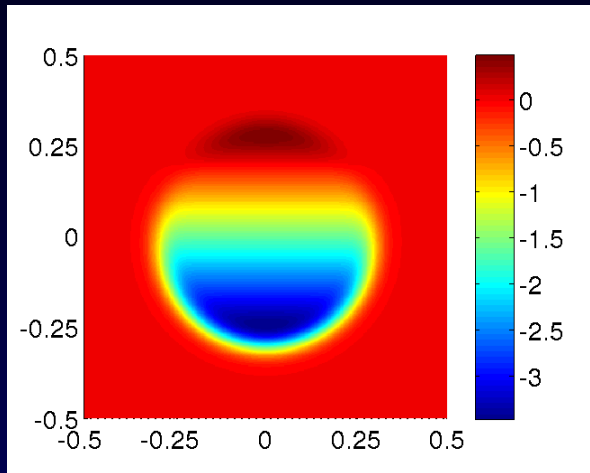
L	N	n	t_{pre}	$t_{\text{new ext}}$	$t_{\text{new int}}$	RAM
4	58081	1120	5 s	0.006 s	0.04 s	230 MB
5	231361	1760	17 s	0.013 s	0.16 s	1 GB
6	923521	3040	78 s	0.05 s	0.7 s	5 GB
7	3690241	5600	6 m	0.16 s	2.7 s	22 GB
8	14753281	10720	22 m	0.5 s	10 s	96 GB

- precomp. time close to $O(N)$, not yet reached $O(N^{3/2})$
- RAM is $O(N \log N)$ ~ 400 complex doubles per unknown

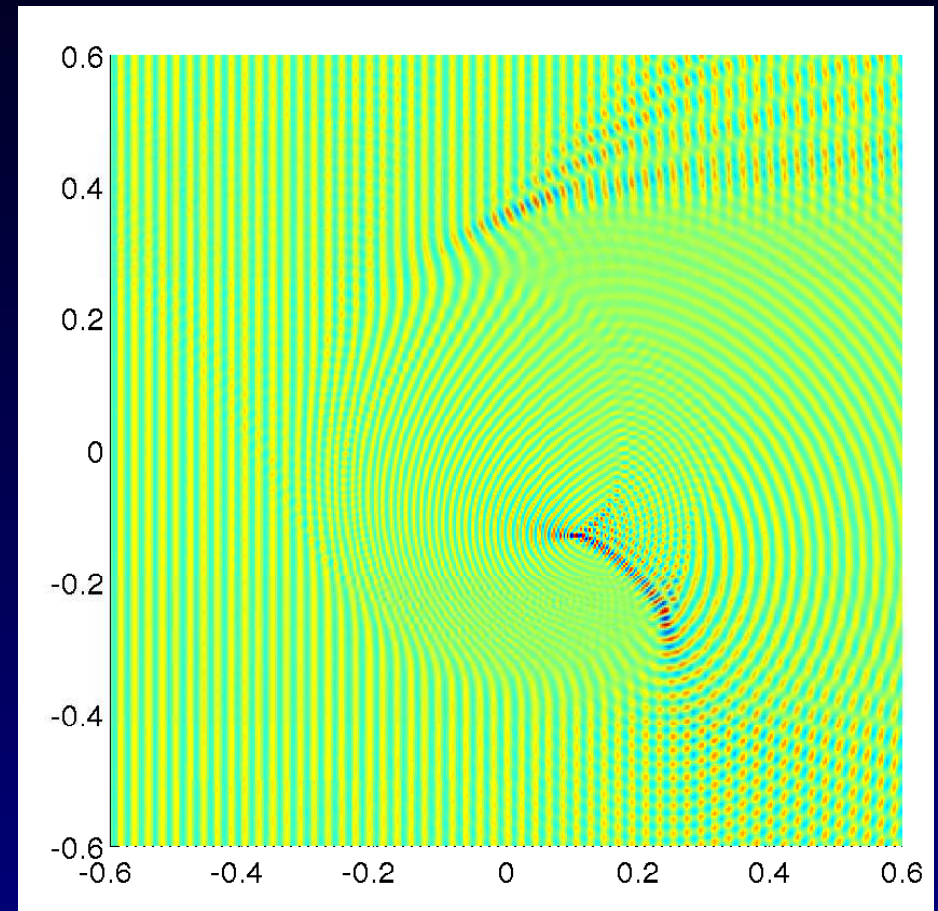
Convergence: graded-index lens, high frequency

plane wave $u^i(\mathbf{x}) = e^{i\kappa \mathbf{d} \cdot \mathbf{x}}$ $\mathbf{d} = (\cos \theta, \sin \theta)$ $\kappa = 300$ 100 shortest- λ 's

$b(\mathbf{x})$:



N	ppw	error	t_{pre}
2.3e5	5	3e-3	17 s
9e5	10	2e-7	78 s
3.7e6	20	7e-10	6 m



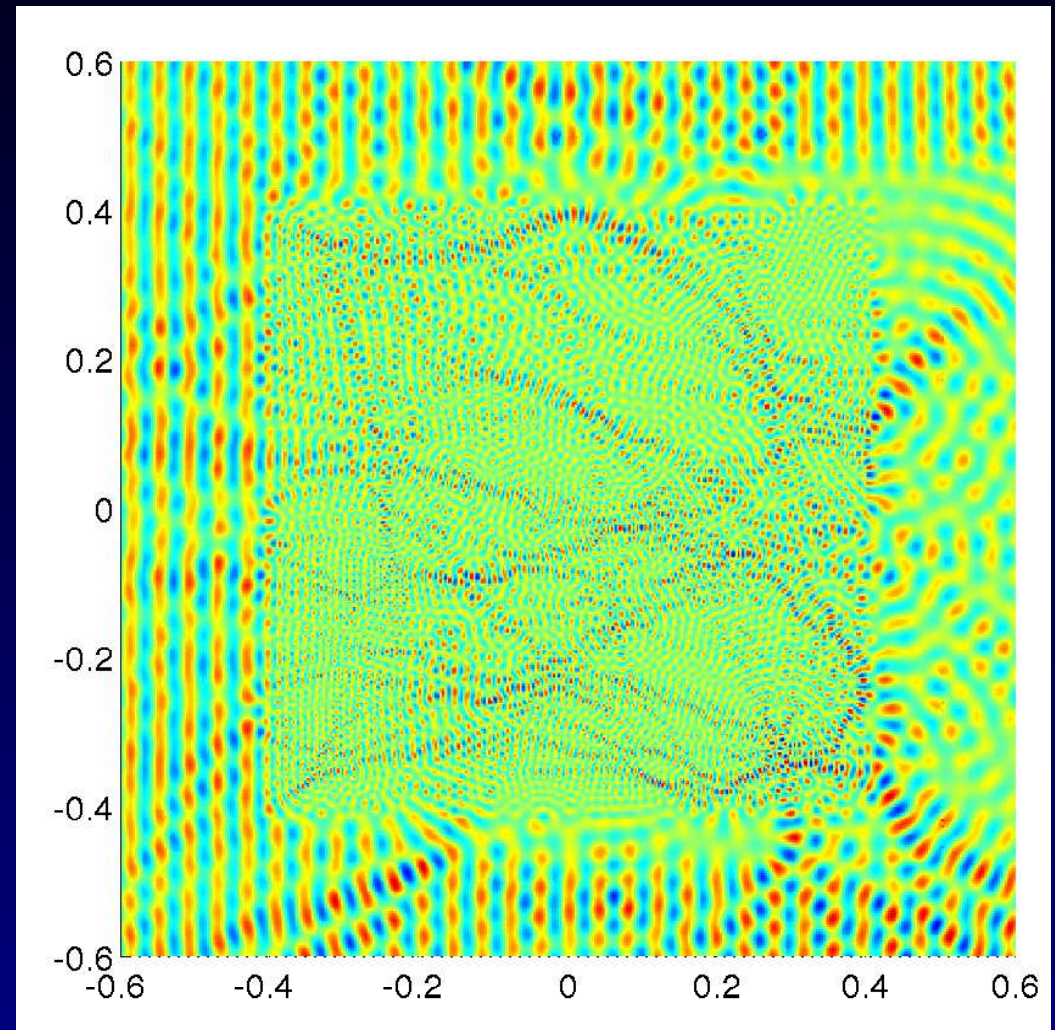
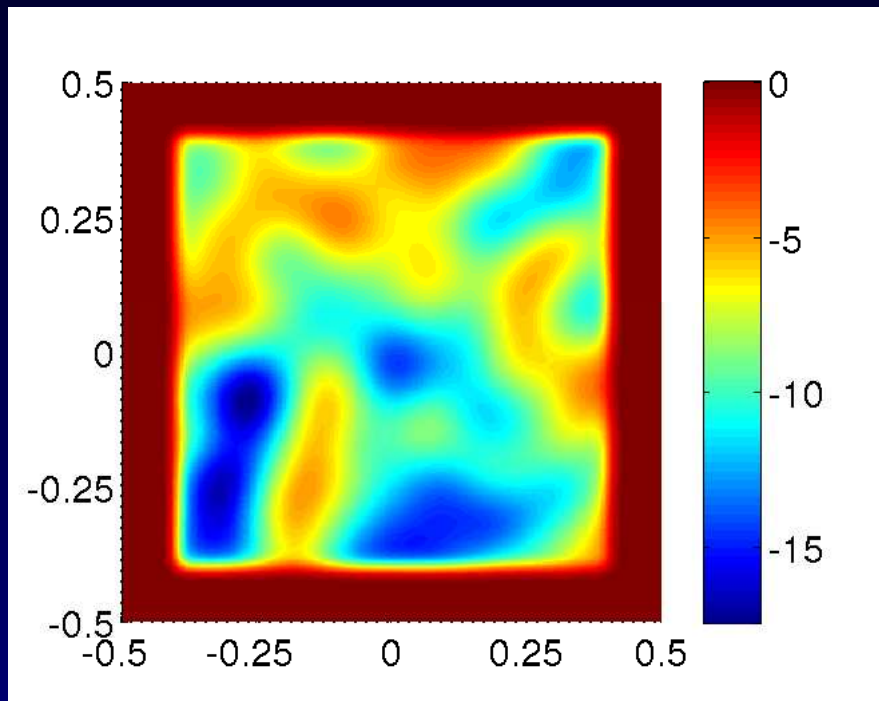
high-order ptwise err rel. to $N=1.4e7$

movie 10^3 angles θ , 3 s each, incl. FMM to 240000 ext. targets

“Bathroom glass” at high frequency

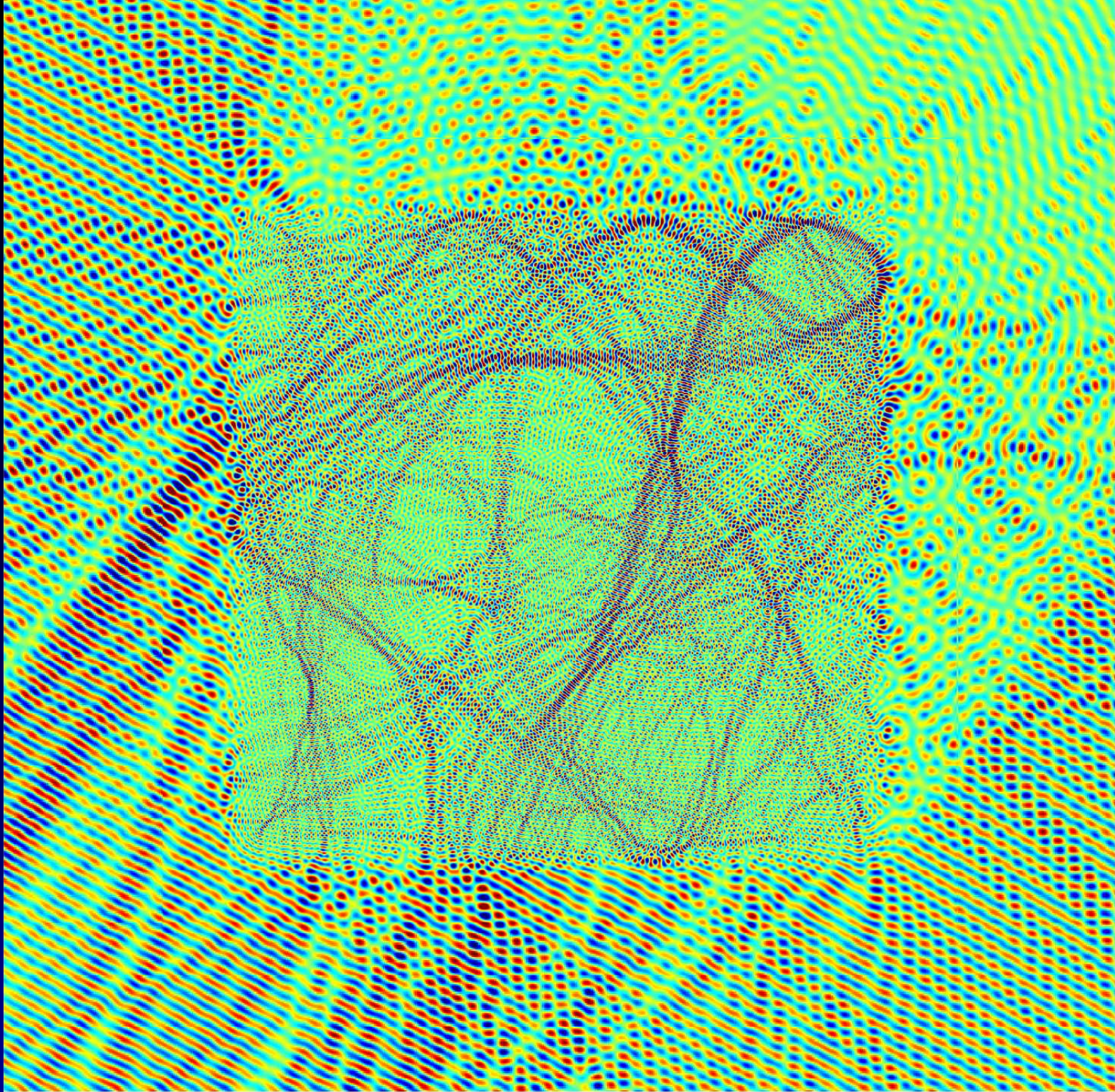
$b(\mathbf{x})$ = random smooth, rolled-off to zero

physics: “2D electron gas”



80 λ on a side time & accuracy similar to before

movie twinkling



$$k = 320$$

200λ across

9 digits (est.)

$$N = 1.4\text{e}7$$

$$t_{\text{pre}} = 22 \text{ min}$$

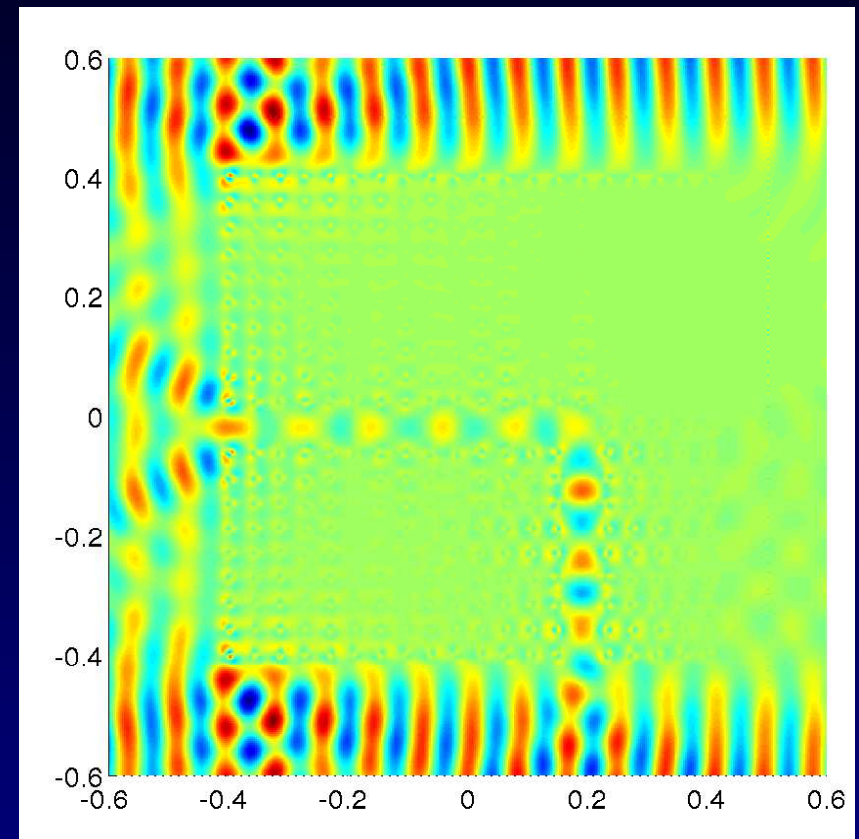
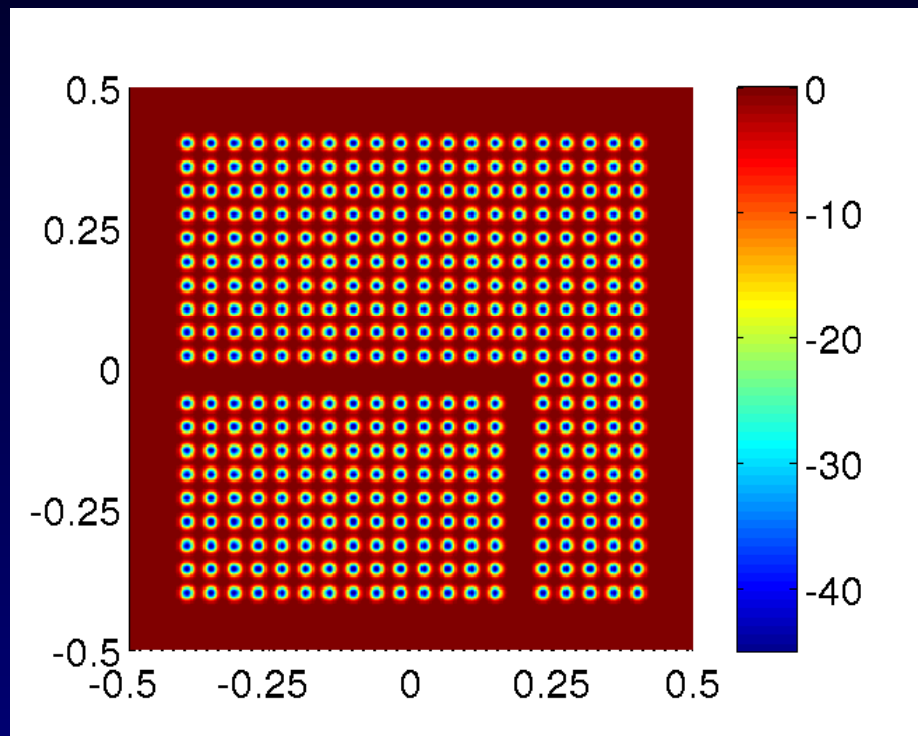
$$t_{\text{new far}} = 0.5 \text{ s}$$

$$t_{\text{new int}} = 10 \text{ s}$$

Photonic crystal waveguide

Until now largely propagating—what if resonances trap waves?

lattice of 400 bumps, each a resonator: choose κ in *bandgap* (can't prop.)



Again, 9 digits $N = 3.7e6$

- also test radially-symmetric $b(\mathbf{x})$ vs sep. of var. soln: stops at 10 digits

Conclusions

$O(N^{3/2})$ direct solver: multiple incident waves at same κ *very fast*
 $N \sim 10^7$, 9 digits @ 200 λ , 100 GB, 20 mins, 0.5 s per u^i , close to $O(N)$

We made robust:

- new proven 2nd-kind BIE
- avoided all box resonances via ItI not DtN

Variants & improvements:

- adaptivity, 3D (Martinsson group); cut leaves, volume source?
- $O(N)$ low- κ DtN version: F^α 's low rank, HBS (Gillman–Martinsson '13)
- $O(N)$ high- κ via butterfly compress F^α 's, R 's ?? $N > 10^8$ 1 TB RAM
- avoid top-level DtN? test my 2nd-kind ItI coupling

Gillman–B–Martinsson, *subm.* BIT '13

Preprints, talks, movies:

funding: NSF DMS-1216656

<http://math.dartmouth.edu/~ahb>