# High-Order Boundary Integral Methods for Slender Bodies

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## **Motivations**



Stokes simulations with fibers are key to modeling complex fluids (suspensions, rheology, industrial, biomedical, cellular biophysics).



Starfish larvae

Drosophila oocyte (Stein et al. 2021)





Mitotic spindle (Nazockdast et al. 2015)

# **Motivations**



Stokes simulations with fibers are key to modeling complex fluids (suspensions, rheology, industrial, biomedical, cellular biophysics).

#### Slender Body Theory (SBT):

- Asymptotic expansion in radius ( $\varepsilon$ ) as  $\varepsilon \to 0$  (Keller-Rubinow '76).
- Doublet correction to make velocity theta-independent (Johnson '80).





Drosophila oocyte (Stein et al. 2021)

Starfish Iarvae (Gilpin et al. 2016)



Mitotic spindle (Nazockdast et al. 2015)

## Slender Body Theory



Error estimates: Rigorous analysis difficult (few very recent studies)

- $\bullet$  classical asymptotics claims:  $\varepsilon^2\log(\varepsilon)$
- rigorous analysis:  $\varepsilon \log^{3/2}(\varepsilon)$  (Mori-Ohm-Spirn '19)
- numerical tests:  $\varepsilon^{1.7}$  (Mitchell et al. '21 -- verify close-touching breakdown) close-to-touching with gap of  $10\varepsilon$ , only 2.5-digits in the infty-norm.  $\varepsilon$ =1e-2 only 1-2 digits achievable by SBT.



Source: http://remf.dartmouth.edu/imagesindex.html

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ε	<b>u</b> <sub>exact</sub>	Rel-Error
le-l	6.1492138359856e-2	0.5e-2
1e-2	9.0984522324584e-2	0.1e-3
1e-3	1.2015655889904e-1	0.2e-5
1e-4	1.4931932907587e-1	0.2e-7
1e-5	1.7848191313097e-1	0.3e-9



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#### Limitations of SBT:

- no convergence analysis for fibers of given nonzero radius.
- uncontrolled errors when fibers close  $O(\varepsilon)$ .

Efficient convergent BIE method needed, allowing adaptivity for close interactions.

## Goals



Solve the slender body BVP

- in a convergent way.
- adaptively when fibers become close.
- efficiently with effort independent of radius.

Validate current SBT simulations.

Focus on rigid fibers in this talk -- flexible fibers for future.

*Related work:* Mitchell et al, '21 (mixed-BVP corresponding to flexible fiber loop)

#### Discretization

#### Geometry description:

- parameterization s along fiber length
- coordinates  $x_c(s)$  of centerline curve
- ullet circular cross-section with radius arepsilon(s)
- ullet orientation vector  $e_1(s)$





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#### Geometry description:

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#### Discretization:

- piecewise Chebyshev (order q) discretization in s for  $x_c(s)$ ,  $\varepsilon(s)$ ,  $e_1(s)$
- Collocation nodes: tensor product of Chebyshev and Fourier discretization in angle with order  $N_{\theta}$ .





#### **Boundary Quadratures**



$$\begin{split} u(x) &= \int_{\Gamma} \mathcal{K}(x-y) \ \sigma(y) \ da(y) \ = \sum_{k=1}^{N_{panel}} \int_{\gamma_k} \mathcal{K}(x-y) \ \sigma(y) \ da(y) \\ &= \underbrace{\sum_{x \notin \mathcal{N}(\gamma_k)} \int_{\gamma_k} \mathcal{K}(x-y) \ \sigma(y) \ da(y)}_{\text{far-field}} \ + \underbrace{\sum_{x \in \mathcal{N}(\gamma_k)} \int_{\gamma_k} \mathcal{K}(x-y) \ \sigma(y) \ da(y)}_{\text{near interactions}} \end{split}$$

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#### Far field approximation:

- Gauss-Legendre quadrature in s.
- periodic trapezoidal rule in  $\theta$ .
- $\bullet$  determine  $\mathcal{N}(\gamma_k)$  using standard error estimates



# **Boundary Quadratures**

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#### Near interactions: for $x \in \mathcal{N}(\gamma_k)$

$$\int_{\gamma_k} \mathcal{K}(x-y) \, \sigma(y) \, da(y) = \int_s \int_{\theta} \mathcal{K}(x-y(s,\theta)) \, \sigma(s,\theta) \, J(s,\theta) \, d\theta \, ds$$

#### Inner integral in $\theta$ :

- potential from a ring source
   -- modal Green's function.
- can be nearly singular as  $x \longrightarrow \gamma_k$ .



#### Outer integral in s:



## Fast Modal Green's Function Evaluation





Analytic representation in special functions - Young, Hao, Martinsson JCP-2012
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- Build special quadrature rules!
  - e.g. generalized Gaussian quadratures: Bremer, Gimbutas and Rokhlin SISC 2010.

# Fast Modal Green's Function Evaluation





- Analytic representation in special functions Young, Hao, Martinsson JCP-2012
   modal Green's functions -- method of choice for axisymmetric problems.
- Build special quadrature rules!
  - e.g. generalized Gaussian quadratures: Bremer, Gimbutas and Rokhlin SISC 2010.
  - Different rule for each nested annular region (up to  $10^{-6}$  from source).
  - $\sim 48$  quadrature nodes for  $n_0=8$  and 10-digits accuracy.
  - $\sim 26M$  modal Green's function evaluations/sec/core (Skylake 2.4GHz)

## Quadratures for Outer Integral



**Near Interactions:** *x* is off-surface or adjacent panel

• panel (Gauss-Lengendre) quadrature with dyadic refienement.



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**Near Interactions:** *x* is off-surface or adjacent panel

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Instead build special quadrature rules!

- replace composite panel quadratures with a single quadrature.
- Separate rules for different aspect ratios (1 --  $10^4$  in powers of 2)

#### Numerical Results - Stokes BVP





**Exterior Stokes** 

 $\begin{array}{ll} \text{Dirichlet BVP:} & \boldsymbol{u}|_{\Gamma} = \boldsymbol{u}_{0}, \\ \Delta \boldsymbol{u} - \nabla p = \boldsymbol{0}, & \boldsymbol{u}(\boldsymbol{x}) \rightarrow \boldsymbol{0} \text{ as } |\boldsymbol{x}| \rightarrow \boldsymbol{0}, \\ \nabla \cdot \boldsymbol{u} = \boldsymbol{0}, \end{array}$ 

wire radius = 1.5e-3 to 4e-3wire length = 16

BIE formulation:  $(\mathcal{I}/2 + \mathcal{D} + \mathcal{S} / (2\varepsilon \log \varepsilon^{-1}))[\boldsymbol{\sigma}] = \boldsymbol{u}_0$ 

Numerical Results - Stokes BVP							FLAT	TIRON TITUTE omputational		
						l leiner				
							1		40	
				1			I-core		40-c	ores
N	N <sub>panel</sub>	$N_{ heta}$	$\epsilon_{\rm GMRES}$	N <sub>iter</sub>	$\left\  e  ight\ _{\infty}$	T <sub>setup</sub> (I	$V/T_{setup}$	$T_{solve}$	40-c T <sub>setup</sub>	ores T <sub>solve</sub>
N 1.5e4	N <sub>panel</sub> 122	Ν <sub>θ</sub> 4	$\epsilon_{_{GMRES}}$ 1e-03	N <sub>iter</sub> 10	$\frac{\ e\ _{\infty}}{1.9\text{e-02}}$	<i>T<sub>setup</sub></i> (1 0.33	$\frac{V/T_{setup}}{(4.4e4)}$	T <sub>solve</sub> 0.7	40-c <i>T<sub>setup</sub></i> 0.024	ores $\frac{T_{solve}}{0.05}$
N 1.5e4 9.1e4	N <sub>panel</sub> 122 252	Ν <sub>θ</sub> 4 12	€ <sub>GMRES</sub> 1e-03 1e-05	<i>N<sub>iter</sub></i> 10 21	$\ e\ _{\infty}$ 1.9e-02 1.7e-04	<i>T<sub>setup</sub></i> ( <i>I</i> 0.33 3.31	$\frac{V/T_{setup}}{(4.4e4)}$ (2.7e4)	<i>T<sub>solve</sub></i> 0.7 61.2	40-c <i>T<sub>setup</sub></i> 0.024 0.197	ores $T_{solve}$ 0.05 5.25
N 1.5e4 9.1e4 9.4e4	N <sub>panel</sub> 122 252 262	Ν <sub>θ</sub> 4 12 12	€ <sub>GMRES</sub> 1e-03 1e-05 1e-07	N <sub>iter</sub> 10 21 33	<i>e</i>    <sub>∞</sub> 1.9e-02 1.7e-04 4.1e-06	<i>T<sub>setup</sub></i> ( <i>I</i> 0.33 3.31 4.43	$\frac{V/T_{setup}}{(4.4e4)}$ (2.7e4) (2.1e4)	<i>T<sub>solve</sub></i> 0.7 61.2 104.3	40-c <i>T<sub>setup</sub></i> 0.024 0.197 0.224	ores <u>T<sub>solve</sub></u> 0.05 5.25 7.69
N 1.5e4 9.1e4 9.4e4 2.0e5	N <sub>panel</sub> 122 252 262 272	N <sub>θ</sub> 4 12 12 24	€ <sub>GMRES</sub> 1e-03 1e-05 1e-07 1e-09	N <sub>iter</sub> 10 21 33 43	<i>∥e∥</i> ∞ 1.9e-02 1.7e-04 4.1e-06 1.4e-08	<i>T<sub>setup</sub></i> ( <i>1</i> 0.33 3.31 4.43 17.70	$\frac{V/T_{setup}}{(4.4e4)}$ (2.7e4) (2.1e4) (1.1e4)	<i>T<sub>solve</sub></i> 0.7 61.2 104.3 586.0	40-c <i>T<sub>setup</sub></i> 0.024 0.197 0.224 0.796	ores <u>T<sub>solve</sub></u> 0.05 5.25 7.69 22.94

#### Numerical Results - close-to-touching





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		l-core		40-cores			
N	$\epsilon_{\rm GMRES}$	N <sub>iter</sub>	$\left\  e  ight\ _\infty$	$T_{setup}$ $(N/T_{setup})$	$T_{solve}$	T <sub>setup</sub>	$T_{solve}$
6.5e4	1e-02	4	2.1e-02	8.1 (8.0e+3)	6.5	1.28	1.4
6.5e4	1e-05	24	2.4e-03	16.8 (3.8e+3)	42.9	2.50	7.7
6.5e4	1e-07	43	2.8e-06	23.5 (2.7e+3)	81.6	3.31	12.8
6.5e4	1e-10	59	5.4e-08	35.6 (1.8e+3)	122.9	4.06	19.2
6.5e4	1e-13	72	1.3e-10	49.9 (1.3e+3)	162.6	5.27	23.2

# Mobility problem

• *n* rigid bodies  $\Omega = \sum_{i=1}^{n} \Omega_i$ with velocities  $\boldsymbol{V}(\boldsymbol{x}) = \boldsymbol{v}_i + \boldsymbol{\omega}_i \times (\boldsymbol{x} - \boldsymbol{x}_i^c)$ , and given forces  $\boldsymbol{F}_i$ , torques  $\boldsymbol{T}_i$  abount  $\boldsymbol{x}_i^c$ .

- Stokesian fluid in  $\mathbb{R}^3 \setminus \Omega$   $\Delta \boldsymbol{u} - \nabla p = 0, \ \nabla \cdot \boldsymbol{u} = 0,$  $\boldsymbol{u} \to 0 \text{ as } \boldsymbol{x} \to \infty.$
- Boundary conditions on  $\partial \Omega$ ,

$$u = V + u_s$$





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- Boundary conditions on  $\partial\Omega$ ,

 $u = V + u_s$ 

unknown:  $oldsymbol{V}(oldsymbol{u}_i,oldsymbol{\omega}_i)$ 





## Mobility problem - double-layer formulation



Represent fluid velocity:  $\boldsymbol{u} = S[\boldsymbol{\nu}(\boldsymbol{F}_i, \boldsymbol{T}_i)] + \mathcal{D}[\boldsymbol{\sigma}]$ and rigid body velocity:  $\boldsymbol{V} = -\sum_{i=1}^{6n} \mathfrak{v}_i \mathfrak{v}_i^T \boldsymbol{\sigma}$ 

Applying boundary conditions ( $\boldsymbol{u} = \boldsymbol{V} + \boldsymbol{u}_s$  on  $\partial \Omega$ ),

$$(\mathcal{I}/2+\mathcal{D})[\boldsymbol{\sigma}]+\sum_{i=1}^{6n}\mathfrak{v}_{i}\mathfrak{v}_{i}^{T}\boldsymbol{\sigma}=\boldsymbol{u}_{s}-\mathcal{S}[\boldsymbol{
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(Pozrikidis - Boundary Integral and Singularity Methods for Linearized Viscous Flow)

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Second kind integral equation, should be well-conditioned. What can possibly go wrong?

# Conditioning of layer-potential operators





• For infinite cylinder (Laplace case):  $\kappa(\mathcal{I}/2 + \mathcal{D}) \sim \varepsilon^{-2} \log^{-1} \varepsilon^{-1}$ 

• Combined field operator well-conditioned:  $\mathcal{I}/2 + \mathcal{D} + \mathcal{S} / (2 \varepsilon \log \varepsilon^{-1})$ 

## Mobility problem - combined field formulation



Represent fluid velocity:  $\boldsymbol{u} = \mathcal{S}[\boldsymbol{\nu}(\boldsymbol{F}_i, \boldsymbol{T}_i)] + \mathcal{K}[\boldsymbol{\sigma}]$ 

and rigid body velocity:  $\mathbf{V} = -\sum_{i=1}^{6n} \mathfrak{v}_i \mathfrak{v}_i^T \boldsymbol{\sigma}$ 

where,  $\mathcal{K} = \mathcal{D} + \mathcal{S}/(2\varepsilon\log\varepsilon^{-1}).$ 

# Mobility problem - combined field formulation



Represent fluid velocity: 
$$\boldsymbol{u} = S[\boldsymbol{\nu}(\boldsymbol{F}_i, \boldsymbol{T}_i)] + \mathcal{K}[\boldsymbol{\sigma} - \sum_{i=1}^{6n} \boldsymbol{v}_i \boldsymbol{v}_i^T \boldsymbol{\sigma}]$$
  
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6n

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where,  $\mathcal{K} = \mathcal{D} + S/(2\varepsilon \log \varepsilon^{-1})$ .

Applying boundary conditions,

$$(\mathcal{I}/2+\mathcal{K})[\boldsymbol{\sigma}-\sum_{i=1}^{6n}\mathfrak{v}_i\mathfrak{v}_i^T\boldsymbol{\sigma}]+\sum_{i=1}^{6n}\mathfrak{v}_i\mathfrak{v}_i^T\boldsymbol{\sigma}=\boldsymbol{u}_s-\mathcal{S}[\boldsymbol{\nu}]$$

Second kind integral equation and well-conditioned!

Time-stepping: 5-th order adaptive SDC

**8-digits accuracy** in quadratures, GMRES solve, and time-stepping.

40 CPU cores





Time-stepping: 5-th order adaptive SDC

**8-digits accuracy** in quadratures, GMRES solve, and time-stepping.

# $\bigcirc\bigcirc$

#### 40 CPU cores





#### Numerical Results - Sedimentation Flow $\cdot 10^{4}$ no-preconditioner --- N 6 100 block-preconditioner 80 4 60 $N_{iter}$ $\geq$ 40 2 20 <sup>0</sup>0 100 200 25Ŏ 50 150 Т

**Close-to-touching:** smaller time-steps, more unknowns (N), high GMRES iteration count (block preconditioner doesn't help).  $\sim 125 \times$  more expensive!



#### 5-th order adaptive SDC

8-digits accuracy in quadratures, GMRES solve, and time-stepping.

# 0.5 million unknowns

64 rings.

160 CPU cores













#### Conclusions



- Convergent boundary integral formulation for slender bodies.
  - unlike SBT, boundary conditions are actually enforced to high accuracy.
- Special quadratures efficient for aspect ratios as large as  $10^5$ .
  - quadrature setup rates up to 20,000 unknowns/s/core (comparable to FMM speeds).
- Stokes mobility problem combined field BIE formulation.
  - well-conditioned formulation for slender-body geometries.
  - high-order time stepping (SDC), block-diagonal preconditioner.

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#### Limitations and ongoing work:

- Open problems: collisions, better preconditions.
- Flexible fibers -- applications in biological fluids.